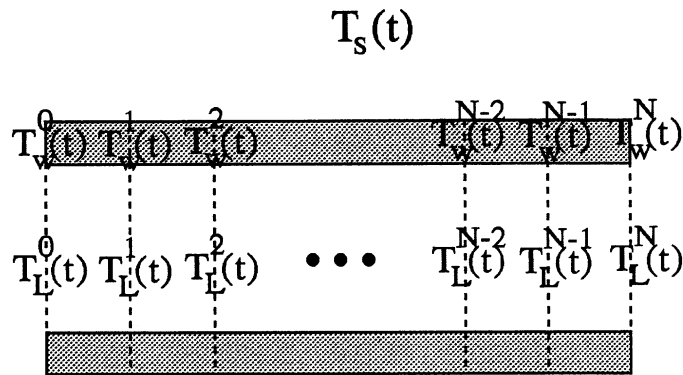


In Finite Difference Method, a grid is established in the direction of spatial variation



$$T_L^j(t) = T_L(z, t)|_{z=j\Delta z}$$

$$T_w^j(t) = T_w(z, t)|_{z=j\Delta z}$$

Now note that each $T_L^j(t)$ and $T_w^j(t)$ depends only on time and therefore are not distributed parameters. In return, we have $2N + 2$ variables in place of the two distributed parameters.

Q: What should be the grid size?

The finer the grid, the more information and the better the approximation. However, it also means more equations.

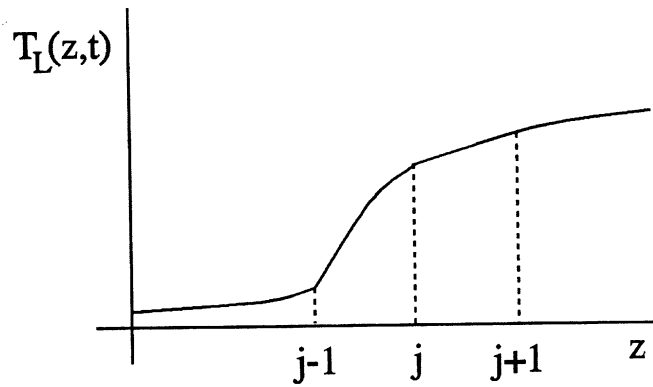
$$\frac{\partial T_L}{\partial t} = -v \frac{\partial T_L}{\partial z} + \frac{1}{\tau_{HL}}(T_w - T_L)$$

$$\frac{\partial T_w}{\partial t} = \frac{1}{\tau_{sw}}(T_s - T_w) + \frac{1}{\tau_{wL}}(T_L - T_w)$$

with $T_s(t)$, $T_L(0, t)$, $T_L(z, 0)$, $T_w(z, 0)$ given, we would like to solve the above at every grid point.

One complication is what to do with the term $\partial T_L / \partial z$.

Finite Difference Approximation



Backward Difference Approximation

$$\left. \frac{\partial T_L}{\partial z} \right|_{z=j\Delta z} \approx \frac{T_L^j(t) - T_L^{j-1}(t)}{\Delta z}$$

Forward Difference Approximation

$$\left. \frac{\partial T_L}{\partial z} \right|_{z=j\Delta z} \approx \frac{T_L^{j+1}(t) - T_L^j(t)}{\Delta z}$$

Average Approximation

$$\begin{aligned} \left. \frac{\partial T_L}{\partial z} \right|_{z=j\Delta z} &\approx \frac{\frac{T_L^j(t) - T_L^{j-1}(t)}{\Delta z} + \frac{T_L^{j+1}(t) - T_L^j(t)}{\Delta z}}{2} \\ &= \frac{T_L^{j+1}(t) - T_L^{j-1}(t)}{2\Delta z} \end{aligned}$$

$$\frac{\partial T_L(z, t)}{\partial t} = -v \frac{\partial T_L(z, t)}{\partial z} + \frac{1}{\tau_{HL}} (T_w(z, t) - T_L(z, t))$$

$$\frac{\partial T_w(z, t)}{\partial t} = \frac{1}{\tau_{sw}} (T_s(t) - T_w(z, t)) + \frac{1}{\tau_{wL}} (T_L(z, t) - T_w(z, t))$$

↓ FDM

At $z = j\Delta z$ (j^{th} grid point)

$$\frac{dT_L^j(t)}{dt} = -v \frac{T_L^{j+1}(t) - T_L^{j-1}(t)}{2\Delta z} + \frac{1}{\tau_{HL}} (T_w^j(t) - T_L^j(t))$$

$$\frac{dT_w^j(t)}{dt} = \frac{1}{\tau_{sw}} (T_s(t) - T_w^j(t)) + \frac{1}{\tau_{wL}} (T_L^j(t) - T_w^j(t))$$

Note that we converted 2 PDEs into 2N ODEs:

$$\frac{dT_L^1(t)}{dt} = -v \frac{T_L^2(t) - T_L^0(t)}{2\Delta z} + \frac{1}{\tau_{HL}} (T_w^1(t) - T_L^1(t))$$

$$\vdots$$

$$\frac{dT_L^N(t)}{dt} = -v \frac{T_L^N(t) - T_L^{N-1}(t)}{\Delta z} + \frac{1}{\tau_{HL}} (T_w^N(t) - T_L^N(t))$$

$$\frac{dT_w^1(t)}{dt} = \frac{1}{\tau_{sw}} (T_s(t) - T_w^1(t)) + \frac{1}{\tau_{wL}} (T_L^1(t) - T_w^1(t))$$

$$\vdots$$

$$\frac{dT_w^N(t)}{dt} = \frac{1}{\tau_{sw}} (T_s(t) - T_w^N(t)) + \frac{1}{\tau_{wL}} (T_L^N(t) - T_w^N(t))$$

Dependent Variables: $T_L^1, \dots, T_L^N, T_w^1, \dots, T_w^N$

Independent Variables: T_s, T_L^0

Comments:

1. For $2N$ ODEs,

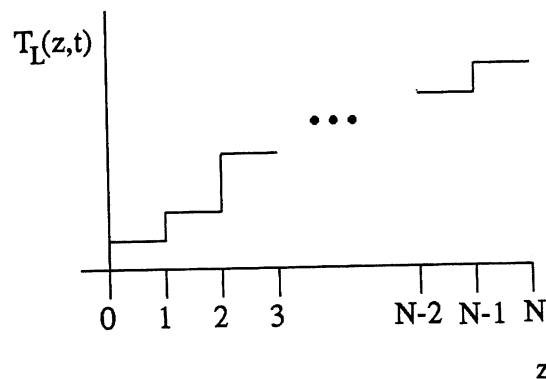
dependent variables are $T_L^1, \dots, T_L^N, T_w^1, \dots, T_w^N$

independent variables are T_s, T_L^0

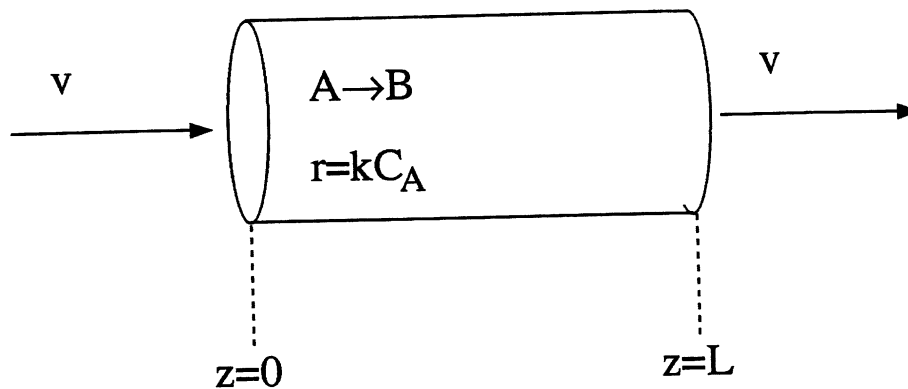
Note that T_L^0 is given by the required boundary condition $T_L(0, t)$ for PDE. So, the boundary condition gets translated to an additional independent variable when discretized.

Of course, to solve the ODEs, the initial conditions $T_L^1(0), \dots, T_L^N(0), T_w^1(0), \dots, T_w^N(0)$ must be given. They come from the initial condition of the PDE $T_L(z, 0)$ and $T_w(z, 0)$.

2. One does not obtain the complete profile, but only values at discrete (grid) points. However, if N is large enough, this will be sufficient. In addition, note that as N increases, the finite difference approximation of the derivatives gets better.
3. Would FDM be equivalent to assuming CST for each segment established by the grids?



C. Isothermal Fixed Bed Reactor



Assumptions:

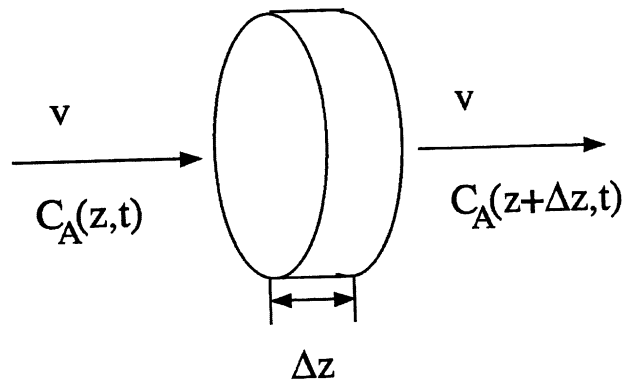
- Isothermal Operation \Rightarrow no need for energy balance
- No radial dispersion \Rightarrow no radial concentration variation
- constant velocity plug flow \Rightarrow no need for total mass balance
- Axial dispersion (due to molecular diffusion and eddy diffusion)

$$\text{Molar Flux} = -\mathcal{D}_L \frac{\partial C_A}{\partial z}$$

Independent Variables: $v, C_A(0, t), C_B(0, t)$

Dependent Variables: $C_A(z, t), C_B(z, t)$

• Shell balance



Component A Balance

$$\frac{d(C_A(z + \Delta z, t)A\Delta z)}{dt} = vAC_A(z, t) - vAC_A(z + \Delta z, t) + \left(-AD_{\mathcal{L}} \frac{\partial C_A}{\partial z'} \Big|_{z'=z} \right) - \left(-AD_{\mathcal{L}} \frac{\partial C_A}{\partial z'} \Big|_{z'=z+\Delta z} \right) - A\Delta z kC_A$$

Divide both sides by $A\Delta z$

$$\frac{dC_A(z + \Delta z, t)}{dt} = -v \frac{C_A(z + \Delta z, t) - C_A(z, t)}{\Delta z} + D_{\mathcal{L}} \frac{\left(\frac{\partial C_A}{\partial z'} \Big|_{z'=z+\Delta z} - \frac{\partial C_A}{\partial z'} \Big|_{z'=z} \right)}{\Delta z} - kC_A$$

Let $\Delta z \rightarrow 0$

$$\boxed{\frac{\partial C_A(z, t)}{\partial t} = -v \frac{\partial C_A(z, t)}{\partial z} + D_{\mathcal{L}} \frac{\partial^2 C_A}{\partial z^2} - kC_A}$$

We obtain a second order partial differential equation. In order to solve the above equation, we need the following conditions.

1. Initial Condition (first order wrt time \Rightarrow one condition)
 $C_A(z, 0)$ (i.e., initial concentration profile across the reactor) must be specified
2. Boundard Condition (second order wrt $z \Rightarrow$ two conditions)
 - $C_A(0, t)$ (i.e., reactor inlet concentration) must be specified
 - $\left. \frac{\partial C_A}{\partial z} \right|_{z=L} = 0$ (no axial dispersion at exit) is another condition typically used.

• Discretization via FDM

Divide the reactor into N segments. Approximate

$$\begin{aligned} \left. \frac{\partial C_A}{\partial z} \right|_{z=j\Delta z} &\approx \frac{C_A^{j+1}(t) - C_A^{j-1}(t)}{2\Delta z} \\ \left. \frac{\partial^2 C_A}{\partial z^2} \right|_{z=j\Delta z} &\approx \frac{\left. \frac{\partial C_A}{\partial z} \right|_{z=(j+\frac{1}{2})\Delta z} - \left. \frac{\partial C_A}{\partial z} \right|_{z=(j-\frac{1}{2})\Delta z}}{\Delta z} \\ &\approx \frac{\frac{C_A^{j+1}(t) - C_A^j(t)}{\Delta z} - \frac{C_A^j(t) - C_A^{j-1}(t)}{\Delta z}}{\Delta z} \\ &= \frac{C_A^{j+1}(t) - 2C_A^j(t) + C_A^{j-1}(t)}{\Delta z^2} \end{aligned}$$

$\frac{\partial C_A}{\partial t} = -v \frac{\partial C_A}{\partial z} + \mathcal{D}_L \frac{\partial^2 C_A}{\partial z^2} - kC_A$ <p style="text-align: center;">↓</p> $\frac{dC_A^j}{dt} = -v \frac{C_A^{j+1} - C_A^{j-1}}{2\Delta z} + \mathcal{D}_L \frac{C_A^{j+1} - 2C_A^j + C_A^{j-1}}{\Delta z^2} - kC_A^j$ <p style="text-align: center;">$j = 1, \dots, N$</p>
--

The final system of equations is N coupled ODEs:

$$\begin{aligned} \frac{dC_A^1}{dt} &= -v \frac{C_A^2 - C_A^0}{2\Delta z} + \mathcal{D}_L \frac{C_A^2 - 2C_A^1 + C_A^0}{\Delta z^2} - kC_A^1 \\ &\vdots \\ \frac{dC_A^N}{dt} &= -v \frac{C_A^{N+1} - C_A^{N-1}}{2\Delta z} + \mathcal{D}_L \frac{C_A^{N+1} - 2C_A^N + C_A^{N-1}}{\Delta z^2} - kC_A^N \end{aligned}$$

Initial Condition

$C_A(z, 0)$ provides $C_A^j(0)$ for $j = 1, \dots, N$

Boundary Conditions

$C_A(0, t)$ provides $C_A^0(t)$ which appears as an independent variable in the above.

$$\left. \frac{\partial C_A}{\partial z} \right|_{z=L} = 0 \Rightarrow \frac{C_A^{N+1} - C_A^{N-1}}{2\Delta z} \approx 0$$

$C_A^{N+1} \approx C_A^{N-1}$. This enables us to get rid of the extra variable C_A^{N+1} that appears in the above.